

# Current Carrying Capacity of Quasi-1D ZrTe<sub>3</sub> Van Der Waals Nanoribbons

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**Abstract**—Quasi-1D van der Waals materials, such as transition metal trichalcogenides, have strong covalent bonds in one direction and weaker bonds in cross-plane directions. They can be prepared as crystalline nanowires or nanoribbons consisting of 1D atomic threads, i.e., chains. We have examined the current carrying capacity of ZrTe<sub>3</sub> nanoribbons using a set of structures fabricated by the shadow mask method. The bulk crystals were synthesized by the chemical vapor transport method and exfoliated onto Si/SiO<sub>2</sub> substrates. It was found that ZrTe<sub>3</sub> nanoribbons reveal an exceptionally high current density, on the order of  $\sim 100$  MA/cm<sup>2</sup>, at the peak of the stressing DC current. The low-frequency noise was of  $1/f$  type near room temperature ( $f$  is the frequency). The noise amplitude scaled with the resistance, following the trend established for other low-dimensional materials. The high current density in ZrTe<sub>3</sub> can be attributed to the single-crystal nature of quasi-1D van der Waals materials.

**Index Terms**—Van der Waals materials, quasi-1D Materials, nanowires; interconnects, current density, ZrTe<sub>3</sub>.

## I. INTRODUCTION

AS AGGRESSIVE scaling in the complementary metal-oxide semiconductor (CMOS) technology continues, there is a growing need to examine new materials that can be used for nanometer-scale local interconnects or device channels [1], [2]. More than a decade ago, the industry moved from aluminum (Al) to copper (Cu) to lower resistance and improve

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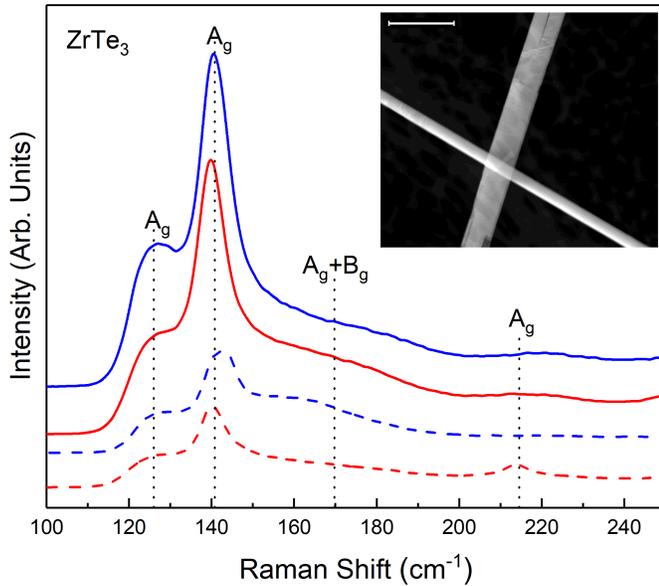
reliability via reduced electromigration. Recently, it has been reported that the 10-nm logic technology node features local interconnects made of cobalt (Co), which enables a 10-fold reduction in electromigration [2]. At present, the current density sustained by Cu interconnects in state-of-the-art CMOS technology is between 2 MA/cm<sup>2</sup> to 3 MA/cm<sup>2</sup> [3], [4].

We recently proposed that quasi-one-dimensional (1D) van der Waals materials, such as transition metal trichalcogenides (TMTs) with formula MX<sub>3</sub> (where M = Nb, Ta, Ti, Zr, and other transition metals; X = S, Se, Te), have properties attractive for applications in nm-scale electronics. In a way similar to transition metal dichalcogenides (TMDs), which exfoliate into 2D layers [5], [6], TMTs exfoliate into quasi-1D atomic thread bundles [7]–[9]. We previously reported that metallic TaSe<sub>3</sub> nanowires have breakdown current density, at the peak of the stressing DC current, on the order of  $\sim 10$  MA/cm<sup>2</sup> [7], [8]. In principle, such quasi-1D materials could be ultimately downscaled by exfoliation, or grown directly, into nanowires with a cross-section of  $\sim 1$  nm  $\times$  1 nm, which corresponds to an individual atomic thread, i.e. chain.

Here we show that nanoribbons made of ZrTe<sub>3</sub>, another member of the TMT family, reveal exceptionally high current density, on the order of  $\sim 100$  MA/cm<sup>2</sup>, at the peak of the stressing DC current (with a cross-section of  $\sim 27$  nm  $\times$  450 nm). Our experimental data also indicate that the low-frequency noise (LFN), another important metric of the material's applicability in electronics, is relatively low, and scales down with the resistance, following the trend found for some other low-dimensional materials.

## II. MATERIAL SYNTHESIS AND CHARACTERIZATION

The selection of ZrTe<sub>3</sub> was based on reports indicating that ZrTe<sub>3</sub> has better electrical conduction properties than other MX<sub>3</sub> materials [10]. For this study, ZrTe<sub>3</sub> crystals were synthesized by the chemical vapor transport (CVT) method [7], [8], [11], [12]. ZrTe<sub>3</sub> crystallizes in the monoclinic space group  $P2_1/m$ . It has strongly anisotropic electronic transport properties, and its bulk electrical resistivity is lower than that of TaSe<sub>3</sub> and several other MX<sub>3</sub> [12]–[15]. However, there are inconsistencies in the literature with respect to the reported electronic band structure of ZrTe<sub>3</sub>, i.e. metallic vs. semiconducting, type of conduction, as well as in the transition temperatures to the charge density wave and superconducting



**Fig. 1.** Raman spectra of  $\text{ZrTe}_3$  for CVT bulk (solid blue) and reference bulk (solid red). The data for exfoliated nanoribbons are shown by dashed lines of the corresponding colors. All spectra were recorded under 633 nm laser excitation and power below 1.5-mW on the surface, to avoid local heating. Inset shows TEM image of the CVT grown exfoliated  $\text{ZrTe}_3$  crystal. The scale bar is  $2 \mu\text{m}$ .

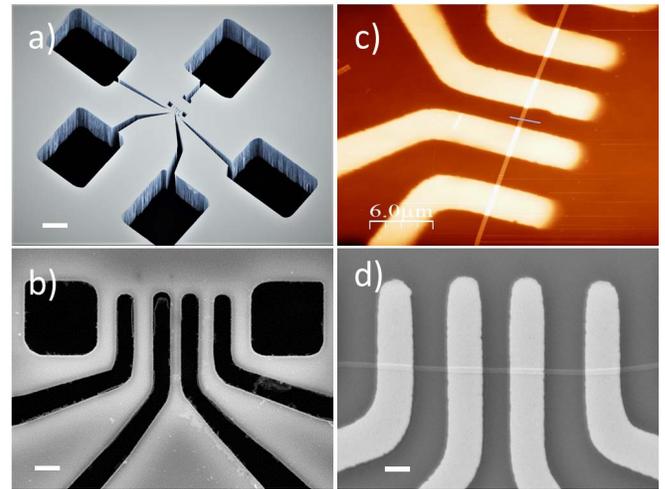
phases [10], [12]–[17]. These issues are not directly related to the present study, which is focused on the current carrying capacity of  $\text{ZrTe}_3$  nanoribbons.

The CVT-grown crystals were characterized by Raman spectroscopy, X-ray diffraction (XRD), scanning electron microscopy (SEM), energy dispersive spectroscopy (EDS), and high-resolution transmission electron microscopy (HR-TEM). **Figure 1** shows Raman spectra (Renishaw) of bulk  $\text{ZrTe}_3$  and exfoliated  $\text{ZrTe}_3$  threads under 633 nm laser excitation for the CVT-grown material and a reference commercial sample (HQ Graphene). All spectra are consistent and in agreement with the literature [18], [19]. As expected for  $P2_1/m$  symmetry,  $\text{ZrTe}_3$  exhibits five main peaks in the range from  $120 \text{ cm}^{-1}$  to  $250 \text{ cm}^{-1}$ .

Mechanical exfoliation onto a  $\text{Si}/\text{SiO}_2$  substrate was used to obtain  $\text{ZrTe}_3$  nanoribbons. One should note that  $\text{ZrTe}_3$  occupies an intermediate position between quasi-2D and quasi-1D materials, owing to its structure and bonding characteristics [12]–[15]. For this reason, it exfoliates into ribbon-like structures with widths larger than thicknesses; thus we refer to nanoribbons of  $\text{ZrTe}_3$  vs. nanowires of  $\text{TaSe}_3$ . The length of the exfoliated nanoribbons could reach hundreds of  $\mu\text{m}$ . While exfoliation can often yield nanoribbons with few-nm lateral dimensions, we focused on the thickness range  $\leq 50 \text{ nm}$  and the width range  $\leq 500 \text{ nm}$ , in order to simplify device fabrication and to ensure the quality of the metal contacts.

### III. DEVICE FABRICATION

We utilized the shadow mask method to directly deposit transmission line measurement (TLM) structures onto pre-selected  $\text{ZrTe}_3$  nanoribbons. The exfoliated  $\text{ZrTe}_3$  nanoribbons were relatively stable in the air (e.g. compared to  $\text{TaSe}_3$ ), which allowed us to use the shadow mask method.

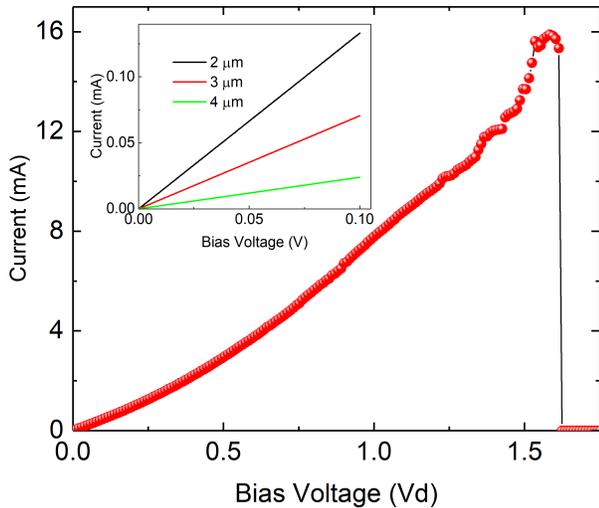


**Fig. 2.** (a) SEM image of a shadow mask with TLM stencil structure patterned on  $500\text{-}\mu\text{m}$  thick  $\text{Si}/\text{SiO}_2$  wafer. (b) Top-view SEM image of the pattern for Ti and Au evaporation to create the source and drain contacts. (c) AFM image of the quasi-1D  $\text{ZrTe}_3$  nanoribbon device, fabricated using the shadow mask. AFM characterization was used to determine the nanowire width and thickness ( $\sim 33\text{-nm}$  in the present case). (d) SEM image of another quasi-1D  $\text{ZrTe}_3$  nanowire device with a different cross-sectional area. The scale bars in (a), (b) and (d) are  $50 \mu\text{m}$ ,  $2 \mu\text{m}$  and  $1 \mu\text{m}$  respectively.

This method was used because it avoids the damage and chemical contamination typically associated with conventional lithographic lift-off processes. It also drastically reduces the total air exposure time. The shadow masks were fabricated using double-side polished Si wafers with  $3 \mu\text{m}$  thermally grown  $\text{SiO}_2$  (Ultrasil Corp.;  $500\text{-}\mu\text{m}$  thickness; P-type;  $<100>$ ). The shadow mask fabrication process began with evaporation of  $200 \text{ nm}$  chromium (Cr) on the front side of the wafer ( $200\text{-nm}$  thickness), followed by a combination of electron beam lithography (EBL) and Cr etchant (1020A) to create a stencil of the TLM pattern. This was followed by fluorine-based reactive ion etching (RIE) to transfer the pattern to the underlying  $\text{SiO}_2$ . Finally, the pattern was transferred into the underlying Si substrate using the deep reactive ion etching (DRIE) (Oxford Cobra). The DRIE etch step was timed to break through to a large backside window that was previously defined using lithographic patterning, RIE, and DRIE. The shadow masks were used to fabricate  $\text{ZrTe}_3$  devices by aligning them with pre-selected nanoribbons on the device substrate, clamping the aligned mask and device substrate together, and placing the clamped assembly in an electron beam evaporator (EBE) for contact deposition ( $10 \text{ nm}$  Ti and  $100 \text{ nm}$  Au) through the shadow mask openings. The completed devices were then transferred to another vacuum chamber for electrical characterization. The thickness and width of  $\text{ZrTe}_3$  nanoribbon were determined from the atomic force microscopy (AFM) inspection. **Figure 2** shows scanning electron microscopy (SEM) images of representative Si shadow masks, as well as AFM and SEM images of the  $\text{ZrTe}_3$  nanoribbon devices fabricated with the masks.

### IV. EXPERIMENTAL RESULTS AND DISCUSSION

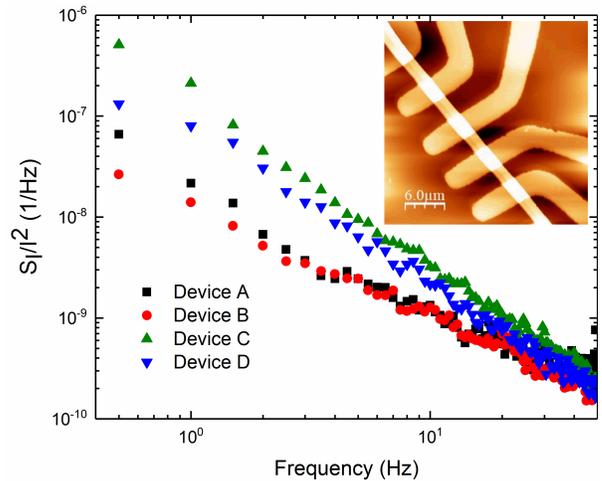
The current-voltage (I–V) characteristics were measured using a probe station (Lakeshore) and a semiconductor



**Fig. 3.** High field I-V characteristics of the best quasi-1D ZrTe<sub>3</sub> nanoribbon device. The apparent breakdown current density, calculated with the AFM measured thickness and SEM measured width, corresponds to  $\sim 10^8$  A/cm<sup>2</sup>, reached at the voltage bias of  $\sim 1.6$  V (with a nanoribbon cross-section of  $\sim 27$  nm  $\times$  450 nm). The current shows some signs of instability at  $V \approx 1.2$  V indicating that some atomic threads started to break. The inset shows low-field I-V characteristics of quasi-1D ZrTe<sub>3</sub> devices with different channel lengths used for the contact resistance extraction. The data indicate the Ohmic nature of the contacts and channel.

analyzer (Agilent B1500). The low-field I-V characteristics of the devices were linear, indicating the Ohmic nature of the contacts and ZrTe<sub>3</sub> channels (see inset to Figure 3). The extracted contact resistance for typical devices was  $2R_C \approx 18\Omega$ , *i.e.* less than 3% of the total ZrTe<sub>3</sub> channel resistance. In Figure 3, we present I-V characteristics of a representative quasi-1D ZrTe<sub>3</sub> nanoribbon device, which was biased up to its complete breakdown. In these devices, the peak of the stressing DC current reached  $\sim 15.8$  mA at the bias voltage of  $\sim 1.6$  V. This corresponds to a maximum breakdown current density on the order of 100 MA/cm<sup>2</sup>. Four other tested devices revealed similar current densities. The electrical resistivity did not show a clear scaling trend, fluctuating in the range from  $\sim 3 \times 10^{-4}\Omega\text{-cm}$  to  $\sim 7 \times 10^{-4}\Omega\text{-cm}$ . The obtained values are in line with the data reported in literature for bulk crystals [14], [15]. One should note that there have been reports suggesting that ZrTe<sub>3</sub> can crystallize in two polymorphs, which have either metallic or semiconducting behavior [16]. However, conditions for phase pure growth have not yet been established, in part because the structural differences between the proposed polymorphs are subtle, making the polymorphs difficult to distinguish by typical characterization techniques (*e.g.*, powder XRD). It is likely that CVT conditions yield a mixture of polymorphs, and thus one can expect a variation in resistivity depending on the dominant phase of the atomic threads.

The breakdown current density obtained for ZrTe<sub>3</sub> is a factor  $\times 10$  larger than the current density at the maximum stressing current measured for Cu wires, tested before Cu introduction in the interconnect technology [20]. It is also a factor of  $\times 3$  larger than the current density at the maximum stressing current for the best TaSe<sub>3</sub> nanowire devices reported to date [7]. Overall, the current density achieved in ZrTe<sub>3</sub>



**Fig. 4.** Normalized noise spectrum density as a function of frequency for ZrTe<sub>3</sub> nanoribbon devices with a cross-section ranging from  $\sim 27$  nm  $\times$  450 nm to  $\sim 100$  nm  $\times$  2000 nm. The inset shows AFM image of one of the fabricated device structures.

nanoribbons is extremely high. This can be attributed to the specific single crystalline structure of quasi-1D van der Waals materials, which minimizes electron scattering at grain boundaries and by interface dangling bonds. Measurements of the low-frequency noise (LFN) are often used to assess the material quality and reliability in conventional devices [21], [22] as well as in 2D materials and devices [23]–[25]. Figure 4 shows the noise characteristics of quasi-1D ZrTe<sub>3</sub> nanoribbons measured at room temperature. LFN is of  $1/f$  type ( $f$  is the frequency). The noise spectral density  $S_I/I^2 \approx 10^{-8} - 10^{-7}\text{Hz}^{-1}$  at  $f = 1$  Hz is comparable to that in some other low-dimensional materials. It scales with resistance,  $R$ , as  $\sim 10^{-11} \times R$ , following the trend observed for carbon nanotubes, graphene and metallic nanowires [8]. The noise level can be expected to decrease further with improved material quality.

## V. CONCLUSIONS

We investigated the current carrying capacity of nanoribbons made from ZrTe<sub>3</sub>, a quasi-1D van der Waals material. The nanoribbon devices revealed an exceptionally high current density, on the order of  $\sim 100$  MA/cm<sup>2</sup>, at the peak of the stressing DC current. Our results suggest that quasi-1D van der Waals materials that consist of single-crystals composed of MX<sub>3</sub> atomic threads have potential for applications in future downscaled electronic technologies.

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